# QED-driven absorption of high-intensity lasers

Tom Blackburn

Chalmers University of Technology

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- How do QED processes affect the absorption of laser light by plasmas?
- When do those processes become significant and how do we study these plasmas?
- What experiments are possible now or in the near future?

## Absorption of laser light Classically

light

laser pulse

matter

electrons

ions

## Absorption of laser light Classically

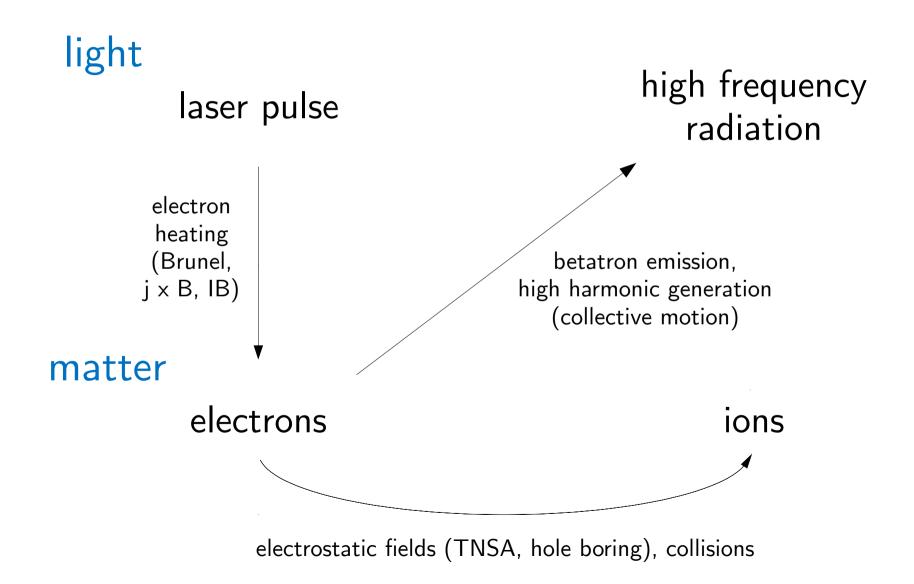
```
light
          laser pulse
        electron
        heating
        (Brunel,
       j x B, IB)
matter
           electrons
```

ions

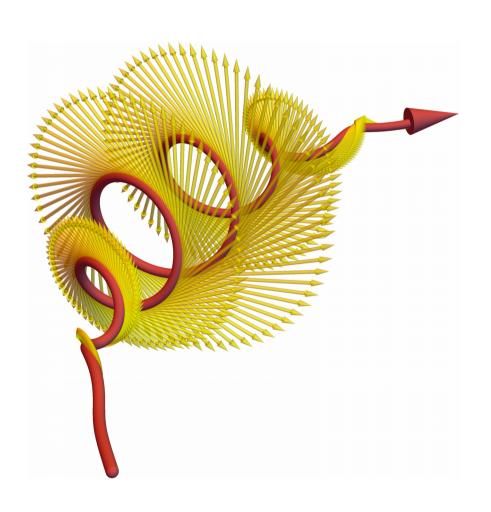
## Classically

```
light
           laser pulse
         electron
         heating
         (Brunel,
        j x B, IB)
matter
            electrons
                                                            ions
                  electrostatic fields (TNSA, hole boring), collisions
```

## Classically



## Radiation from accelerating electrons

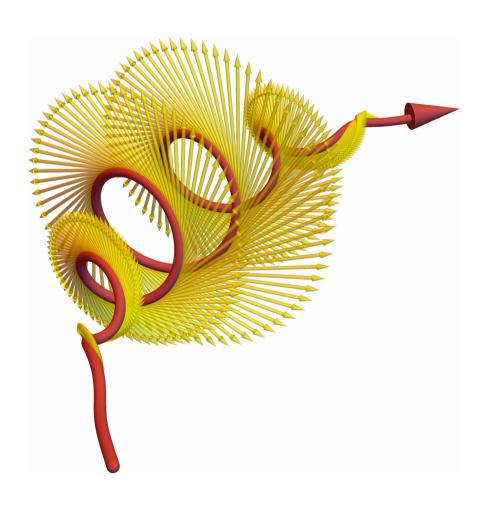


 Fractional energy loss to radiation in one laser period is

$$\frac{P_{\text{rad}}\tau_{\text{L}}}{\gamma mc^{2}} = 0.1 \left(\frac{I}{2 \times 10^{22} \text{ Wcm}^{-2}}\right)^{3/2}$$

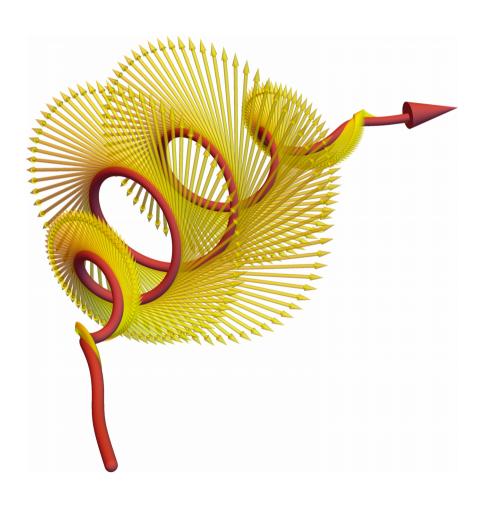
• i.e. 10% at the current intensity frontier.

#### Radiation reaction



- Radiation carries momentum as well as energy, so the electron must recoil.
- The Lorentz force does not account for that recoil.
- Modify electron equation of motion.

## From classical to quantum radiation reaction



- Many different formulations of the classical radiation reaction force<sup>1</sup>.
- Lorentz-Abraham-Dirac and Landau-Lifshitz are the most widely accepted.
- Can be derived<sup>2</sup> from QED, and are consistent to first order in  $\alpha$ .

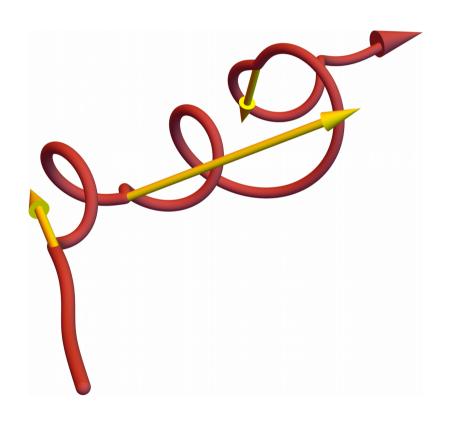
## From classical to quantum radiation reaction



- When does a classical treatment of radiation become insufficient?
- Important parameter is  $\chi$  (ratio of electric field in electron rest frame to the Schwinger field)

$$\chi = rac{\gamma \left| \mathbf{E} + \mathbf{v} imes \mathbf{B} 
ight|}{E_{\mathsf{Sch}}}$$

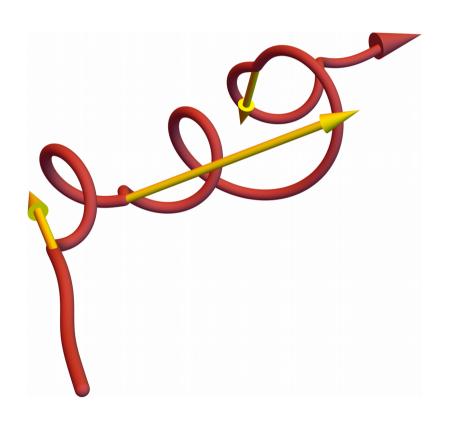
## Gamma ray emission



- A electron in a laser field has Lorentz factor about  $a_0$
- So in a typical laserplasma interaction

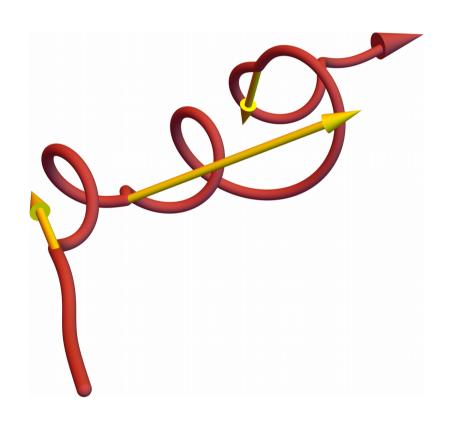
$$\chi = \frac{\mathit{I}\lambda}{\mathsf{5.7}\times\mathsf{10^{23}~Wcm^{-2}}\mu\mathsf{m}}$$

## Gamma ray emission



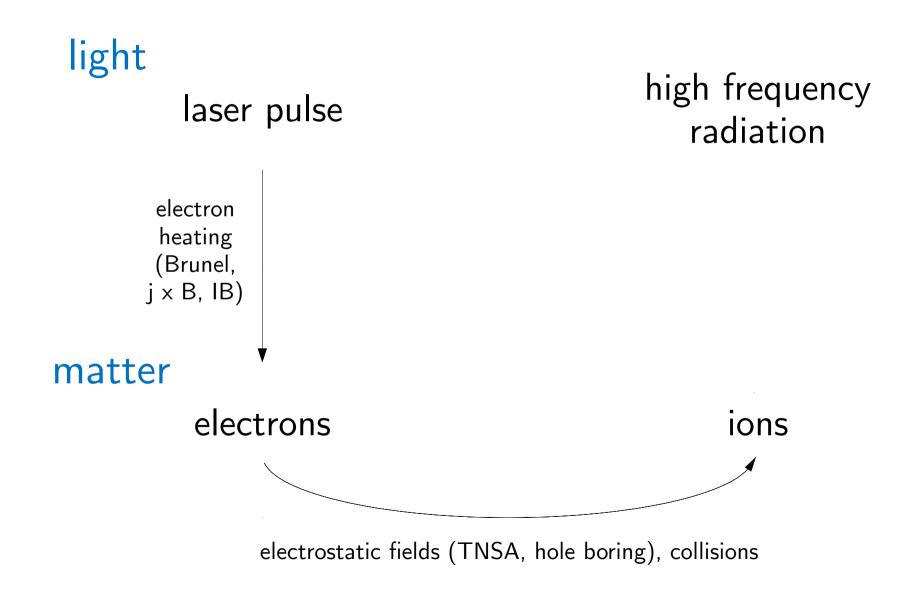
- The typical photon energy as a fraction of the electron energy is  $0.44\chi$ .
- Electrons can lose a substantial fraction of their energy in a single emission.
- Emission becomes a stochastic process.

## Gamma ray emission

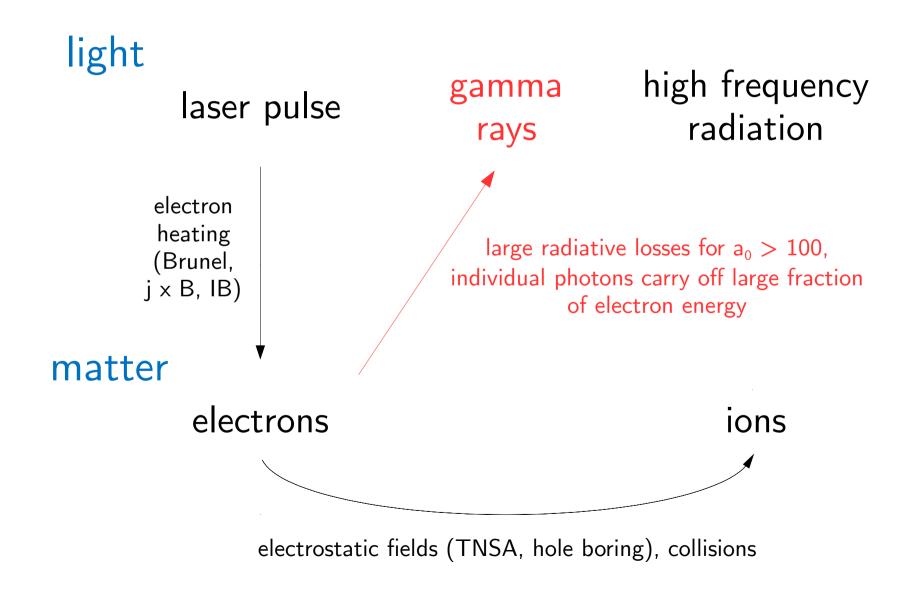


- But even for  $\chi = 0.1$ , quantum corrections to the photon spectrum mean that a classical treatment overestimates the radiated power by 30%.
- The classical force must be damped by a factor  $g(\chi) < 1$ .

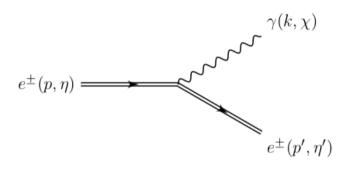
## With gamma ray emission

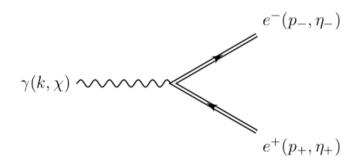


## With gamma ray emission



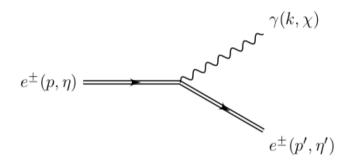
## Electron-positron pair creation

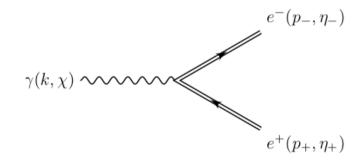




- If we can emit photons with multi-MeV energies, what about pair creation?
- Breit-Wheeler pairs are created when gamma rays collide with photons of the laser pulse.
- Purely quantum phenomenon.

## Electron-positron pair creation

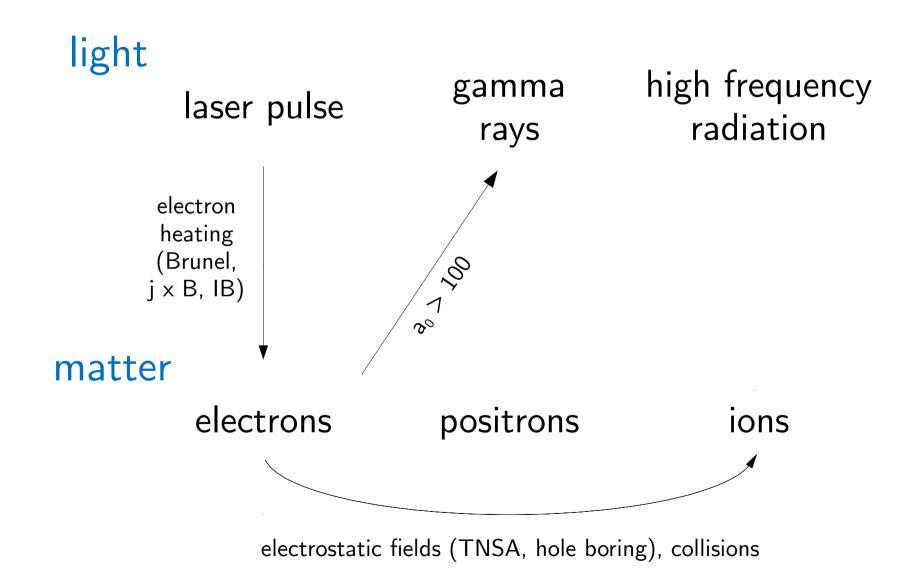




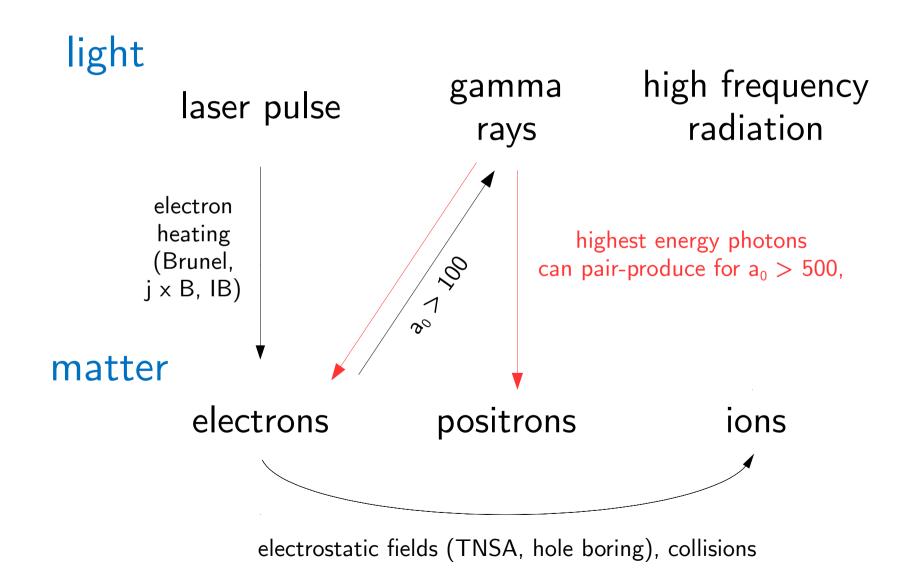
- Pair creation probability depends on the photon χ and the distance over which the intense fields are sustained.
- Result is

$$rac{d au}{dz} \simeq 0.6 \left(rac{I}{10^{22} \ ext{Wcm}^{-2}}
ight)^{1/2} \ \exp\left(-rac{4}{3\chi}
ight) \mu ext{m}^{-1}$$

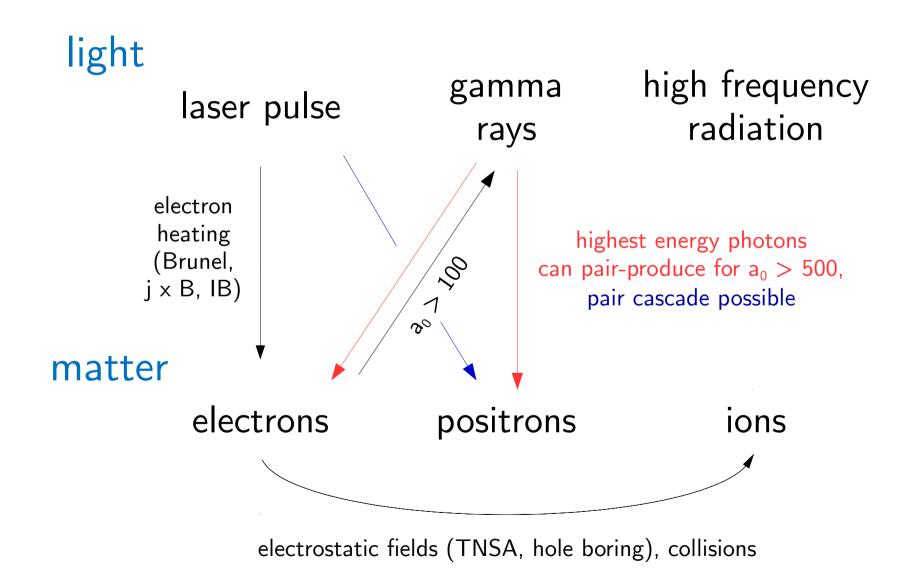
## With pair production



## With pair production

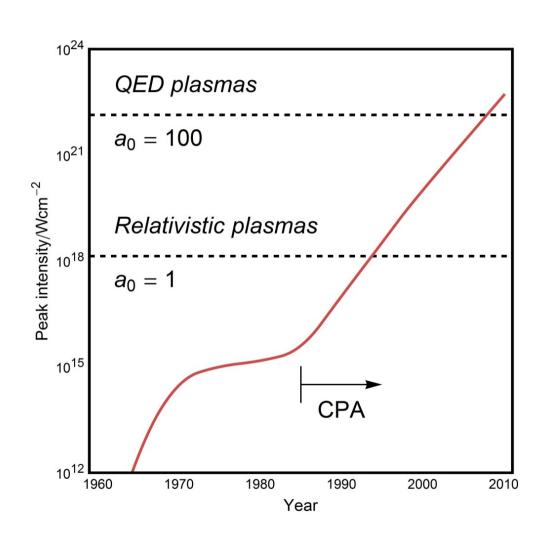


## With pair production



#### Why study QED-plasmas?

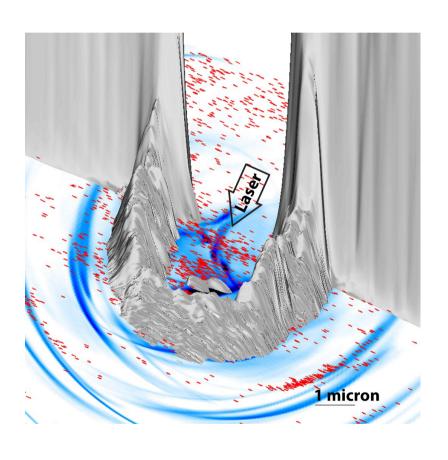
## Increasing intensity of short-pulse lasers



- Current intensity record is  $2 \times 10^{22}$  Wcm<sup>-2</sup> at the HERCULES laser, Michigan.
- Next generation laser facilities (ELI etc.) will exceed 10 PW, 10<sup>23</sup> Wcm<sup>-2</sup>.

#### How can we study QED-plasmas?

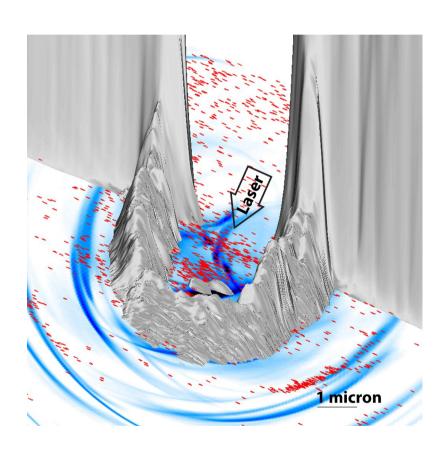
### Laser-solid interactions at >10 PW



- Important physics across a wide range of timescales.
- 10<sup>12</sup> particles in the initial state.
- QED processes in a strong background field with complex structure.

#### How can we study QED-plasmas?

### Laser-solid interactions at >10 PW



- Make simplifying assumptions to study these interactions with PIC codes.
- Monte-Carlo modelling of QED events (Duclous et al, PPCF 53, 015009).
- e.g., EPOCH, OSIRIS, PICADOR...

## QED-PIC codes

- If the formation length is short enough, we can assume the fields are quasi-static over the emission process.
- The formation length is smaller than the laser wavelength by a factor of the strength parameter a<sub>0</sub>.

- We also assume the fields are weak.
- i.e. the fields invariants *f* and *g* are vanishingly small

$$f = |E^2 - B^2|/E_{\text{crit}}^2$$
$$g = |\mathbf{E} \cdot \mathbf{B}|/E_{\text{crit}}^2$$

## QED-PIC codes

- The rates for photon and pair production can then be calculated in an equivalent system of fields with the same instantaneous value of χ.
- e.g. a static magnetic field in the ultrarelativistic limit: T. Erber, Rev. Mod. Phys. 38, 626 (1966)

- Other constraints on PIC simulations:
- Size of timestep must be set to avoid multiple scatterings (constraint set by rate of photon emission is strongest).
- C. P. Ridgers et al, J.Comp. Phys., 260, 273 (2014)

## e.g. laser absorption fraction to gamma-rays

$$f_e = \begin{pmatrix} ext{electron} \\ ext{temperature} \\ ext{a}_0 mc^2 \end{pmatrix} imes \begin{pmatrix} ext{hot electron} \\ ext{fraction} \\ ext{f}_h \end{pmatrix} imes \begin{pmatrix} ext{number} \\ ext{density} \\ ext{a}_0 n_{ ext{cr}} \end{pmatrix} / \begin{pmatrix} ext{laser energy} \\ ext{density} \\ f_2 a_0^2 mc^2 n_{ ext{cr}} \end{pmatrix}$$

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$$f_{\gamma} = \begin{pmatrix} \text{radiated} \\ \text{energy} \\ \propto \chi^2 g(\chi) \end{pmatrix} \times \begin{pmatrix} \text{hot electron} \\ \text{fraction} \\ f_h \end{pmatrix} \times \begin{pmatrix} \text{number} \\ \text{density} \\ a_0 n_{\text{cr}} \end{pmatrix} / \begin{pmatrix} \text{laser energy} \\ \text{density} \\ \frac{1}{2} a_0^2 m c^2 n_{\text{cr}} \end{pmatrix}$$

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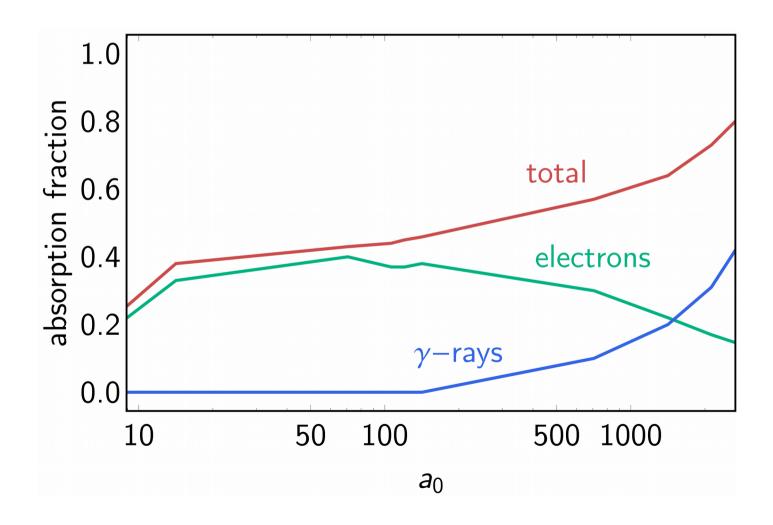
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$$f_{\gamma} = \begin{pmatrix} \text{radiated} \\ \text{energy} \\ \propto a_0^4 g(\chi) \end{pmatrix} \times \begin{pmatrix} \text{hot electron} \\ \text{fraction} \\ f_h \end{pmatrix} \times \begin{pmatrix} \text{number} \\ \text{density} \\ a_0 n_{\text{cr}} \end{pmatrix} / \begin{pmatrix} \text{laser energy} \\ \text{density} \\ \frac{1}{2} a_0^2 m c^2 n_{\text{cr}} \end{pmatrix}$$

$$\propto f_h a_0^3 g(\chi)$$

## e.g. laser absorption fraction to gamma-rays

2D QED-PIC simulations of 40fs of 1 micron light striking a plasma with density just above critical:



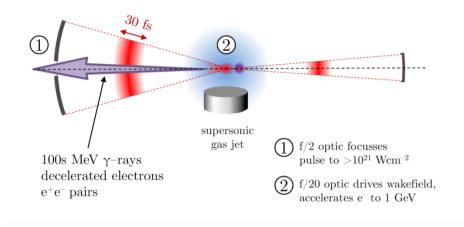
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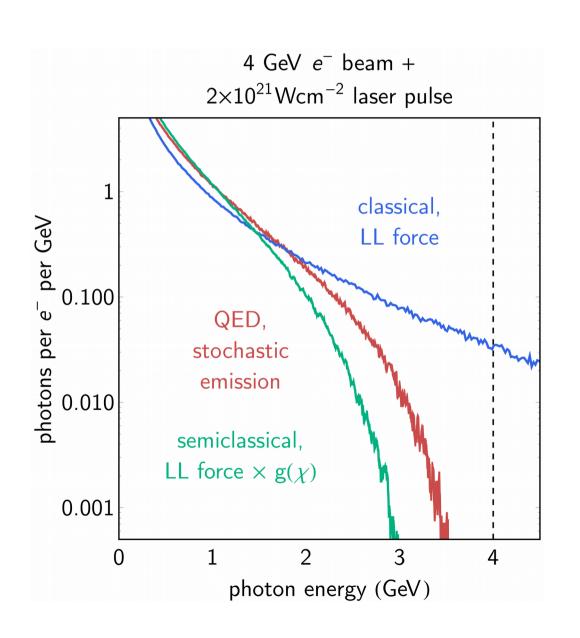
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$$\propto f_h a_0^3 g(\chi)$$

When  $a_0$  is high enough that a substantial fraction of the laser energy is absorbed to gamma rays, we also have  $\chi > 0.1$ .

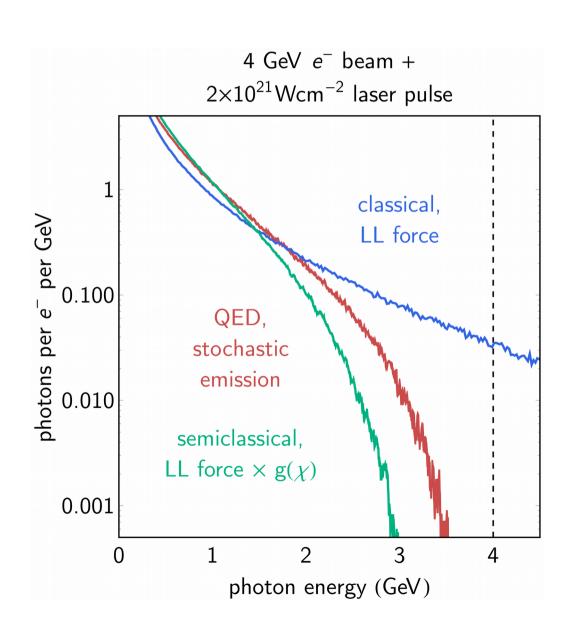


- Accelerate electrons to high energies before they encounter a high intensity field.
- Use gamma ray spectra to distinguish quantum from classical radiation reaction.

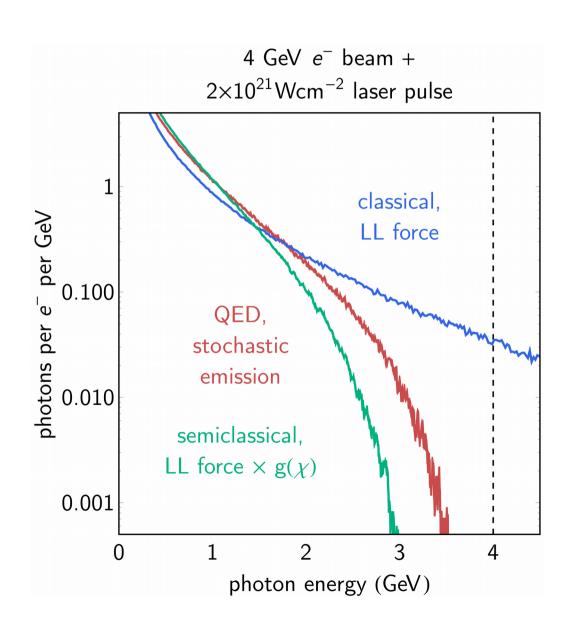


• 
$$\chi = \gamma E / E_{Sch}$$

- Accelerate electrons to high energies before they encounter a high intensity field.
- Use gamma ray spectra to distinguish quantum from classical radiation reaction.



- "Photon energy" in classical emission is unbounded, overestimates high-energy tail (a lot).
- Correcting the radiated intensity by  $g(\chi)$ , but without treating emission probabilistically, underpredicts that tail.



- Reproducibility of electron beam – given the initial energy and the photon spectra, determine a<sub>0</sub> of target laser consistent with either classical and quantum emission.
- Only one of these will be consistent with the actual laser pulse.

- Next generation laser facilities will produce plasmas dominated by radiation reaction and QED processes.
- We can begin to study these interactions using QED-PIC simulations.
- But we can (and need to) to validate the model used in those simulations with colliding-beam experiments possible with the BELLA facility.